

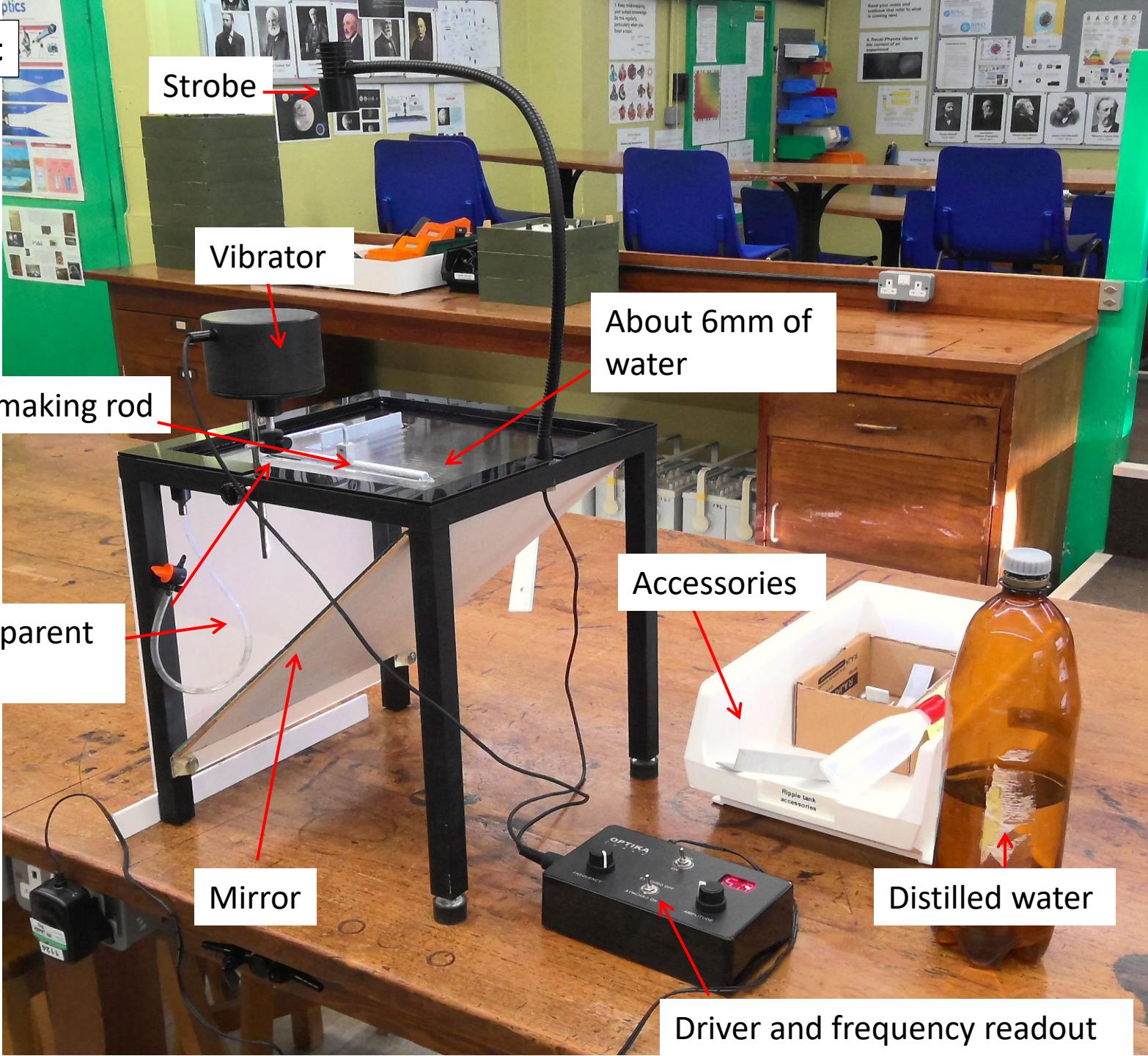
Ripple tank

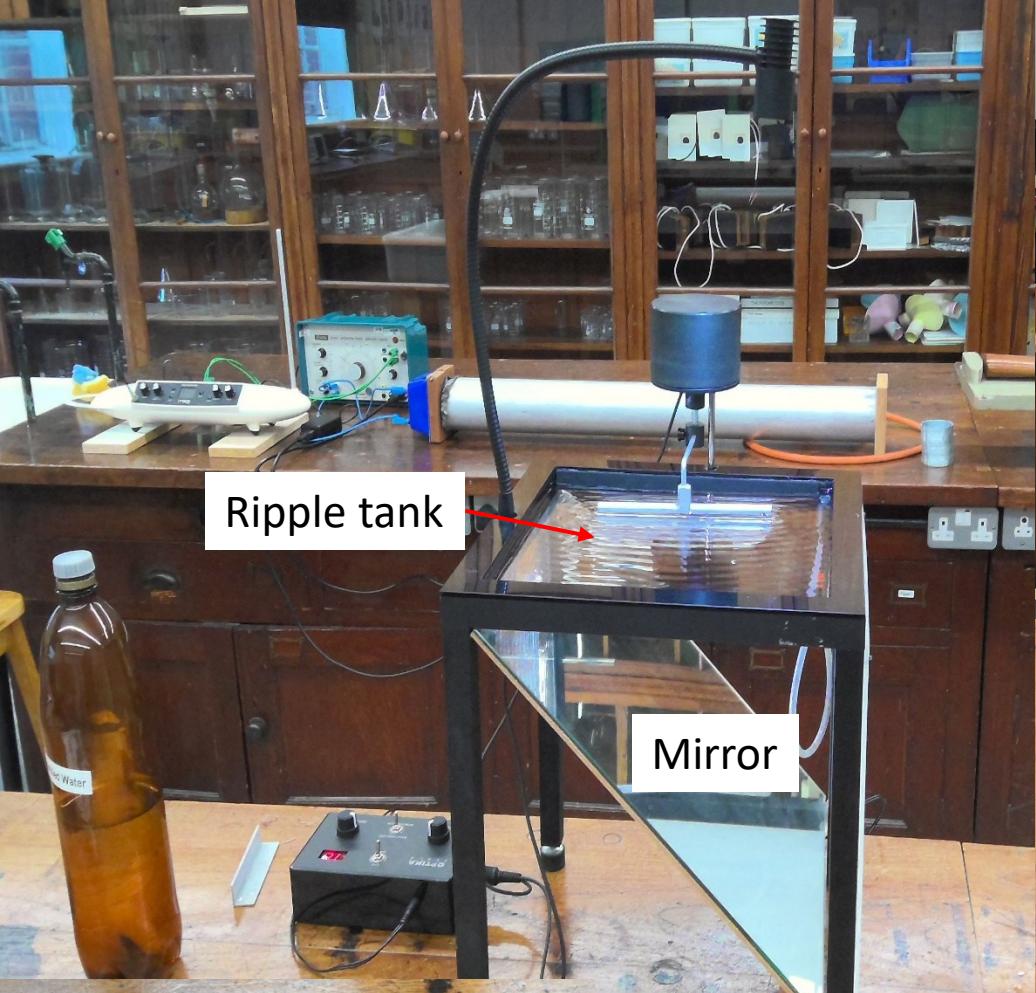
$$\omega = 2\pi f, \quad k = 2\pi/\lambda$$

$$c_p = \frac{\omega}{k} = f\lambda, \quad c_g = \frac{d\omega}{dk}$$

$$\omega^2 = \left(\frac{\sigma k^3}{\rho_1} + gk \right) \tanh(kD)$$

Equipment





Ripple tank

Mirror



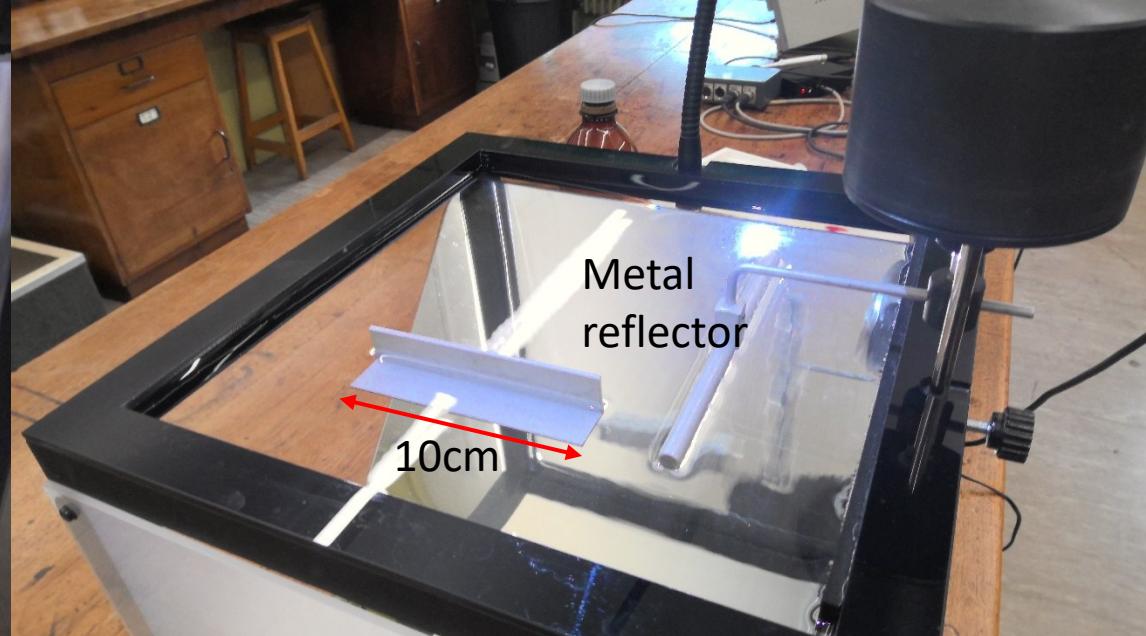
Observe
wavefronts
on screen

Synchro on sets strobe to wave frequency, so waves appear static



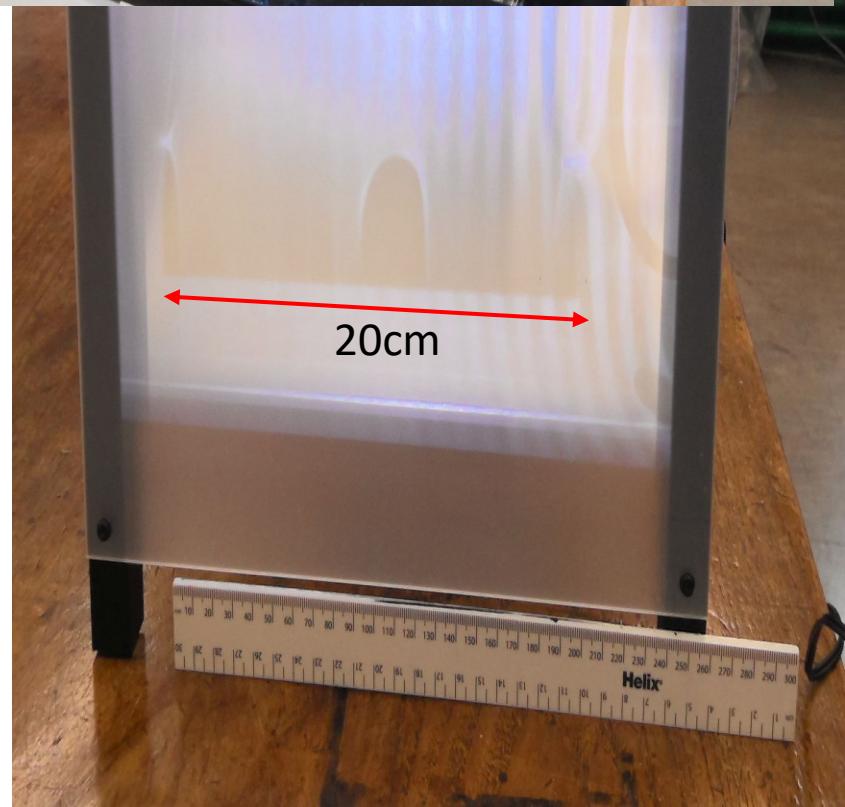
Wave driver control (main powered).
Range of frequencies is 0 to 50Hz. You will probably struggle
to make meaningful wavelength measurements below
about 13 Hz.

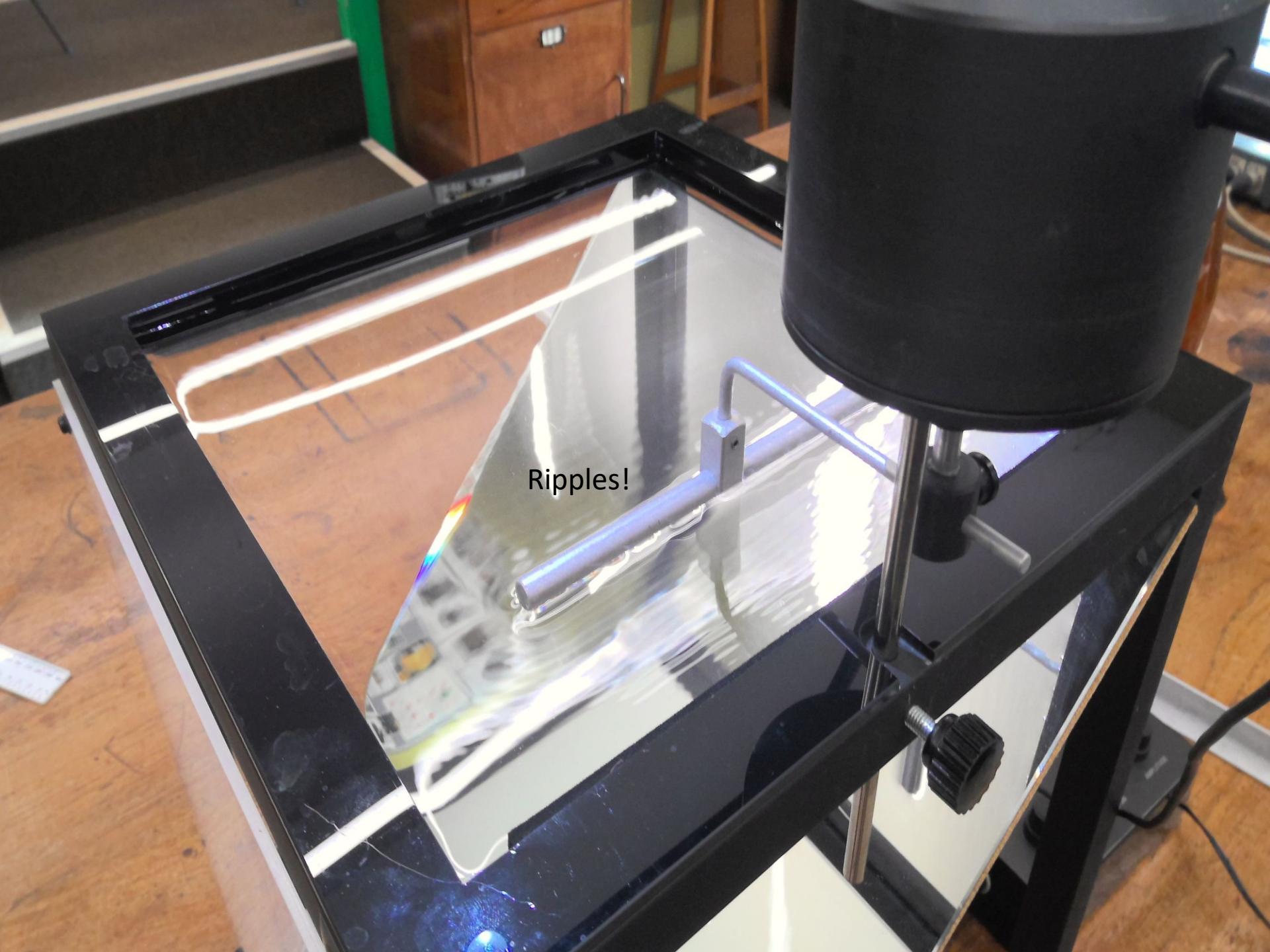
Don't forget to
remove the reflector
from the tank!



Measuring off the semi-transparent screen
is a **magnification factor of x 2.00**.

You can see this by placing one of the (10cm
long) accessories in the tank and measuring the
length of the shadow (20cm).



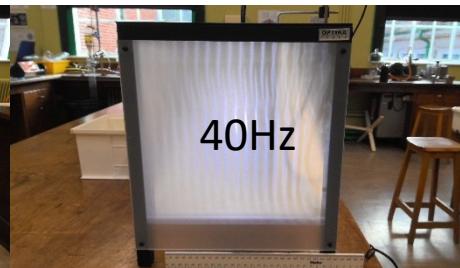


Ripples!

The image shows a physics demonstration setup. A rectangular tank filled with water is mounted on a black metal frame. A cylindrical black container is positioned above the tank, with a blue tube connected to it, likely for dispensing water. The tank has a black frame and a glass front. The water surface is disturbed, creating ripples. A white rectangular tape measure is placed across the top of the tank. The background shows a wooden floor and some wooden furniture, including a cabinet and a chair. The text "Ripples!" is overlaid on the image, pointing to the disturbed water surface.

The photographs below illustrate varying the ripple tank frequency from 5Hz to 50Hz.

I used a water depth of 6mm. You can easily make direct wavelength measurements from the screen using a ruler. **Just remember to halve the measurements due to the magnification!**



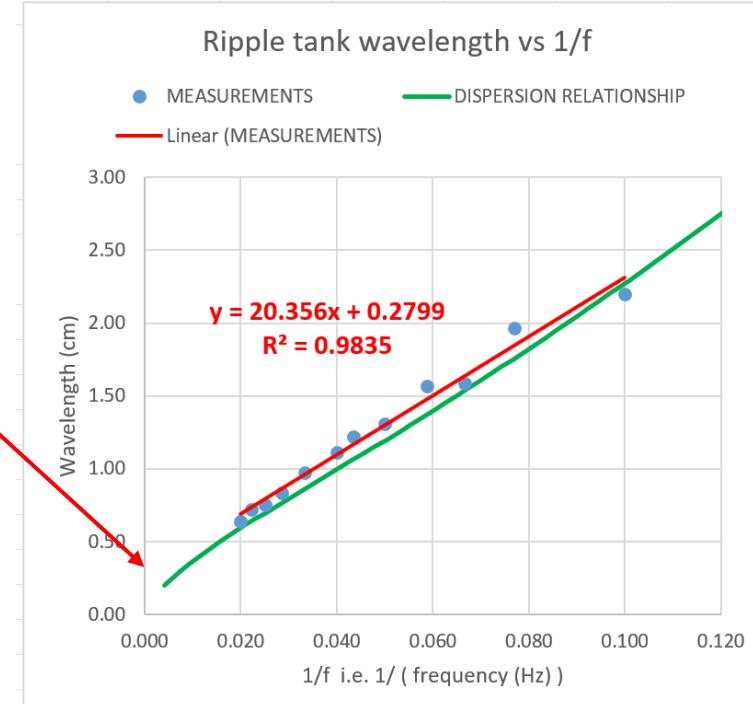
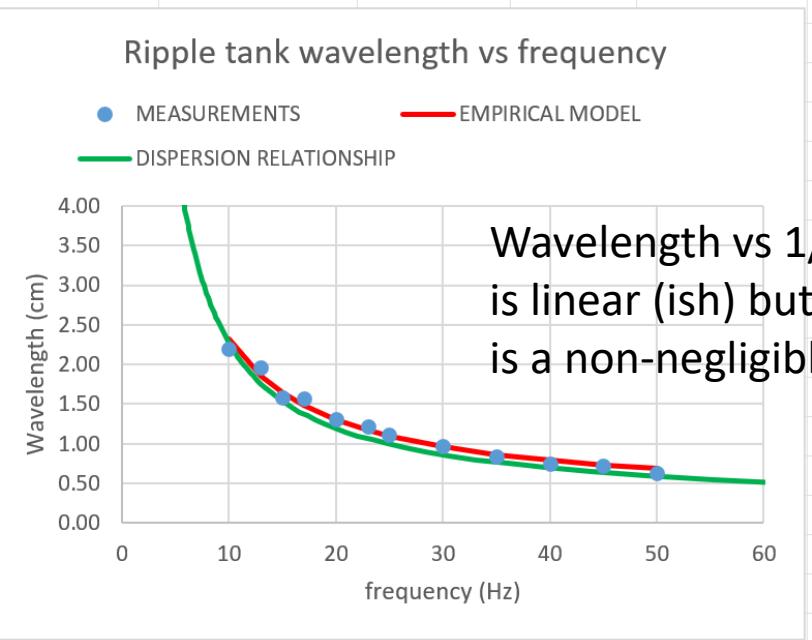
$$c_p = f\lambda$$

Observe that the *phase velocity* of ripples appears to increase from about 22 cm/s to 32 cm/s between 10Hz and 50Hz.

Empirical Model

$$\lambda/\text{cm} \times \frac{20.4}{(f/\text{Hz})} + 0.28$$

from strobe box		measurements from photo					
frequency (Hz)	1/f (Hz ⁻¹)	Pixels for N waves	Pixels for 30 cm ruler	Number of wavelengths N	wavelength /cm	phase velocity $c = f\lambda$ (cm/s)	MODEL wavelength (cm)
10	0.100				2.20	22.0	2.32
13	0.077	1066	1628	5	1.96	25.5	1.85
15	0.067	841	1594	5	1.58	23.7	1.64
17	0.059	1149	1572	7	1.57	26.6	1.48
20	0.050	1042	1494	8	1.31	26.2	1.30
23	0.043	1198	1640	9	1.22	28.0	1.17
25	0.040	893	1509	8	1.11	27.7	1.10
30	0.033	887	1523	9	0.97	29.1	0.96
35	0.029	853	1533	10	0.83	29.2	0.86
40	0.025	800	1594	10	0.75	30.1	0.79
45	0.022	885	1541	12	0.72	32.3	0.73
50	0.020	635	1497	10	0.64	31.8	0.69



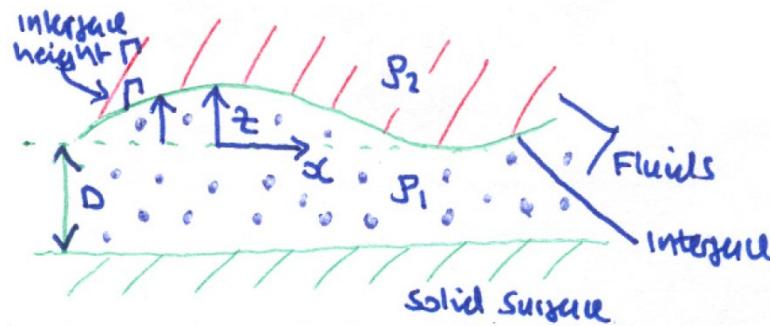
Surface tension of water [edit]

The surface tension of pure liquid water in contact with its vapor has been given by IAPWS^[40] as

$$\sigma = 235.8 \left(1 - \frac{T}{T_C}\right)^{1.256} \left[1 - 0.625 \left(1 - \frac{T}{T_C}\right)\right] \text{ mN/m},$$

where both T and the critical temperature $T_C = 647.096 \text{ K}$ are expressed in [kelvins](#). The region of validity the entire vapor–liquid saturation curve, from the triple point ($0.01 \text{ }^\circ\text{C}$) to the critical point. It also provides reasonable results when extrapolated to metastable (supercooled) conditions, down to at least $-25 \text{ }^\circ\text{C}$. This formulation was originally adopted by IAPWS in 1976 and was adjusted in 1994 to conform to the International Temperature Scale of 1990.

The uncertainty of this formulation is given over the full range of temperature by IAPWS.^[40] For temperatures below $100 \text{ }^\circ\text{C}$, the uncertainty is $\pm 0.5\%$.



Water surface tension (N/m) at 25 deg C 0.0731 Temperature /K 291

strength of gravity (N/kg) 9.8

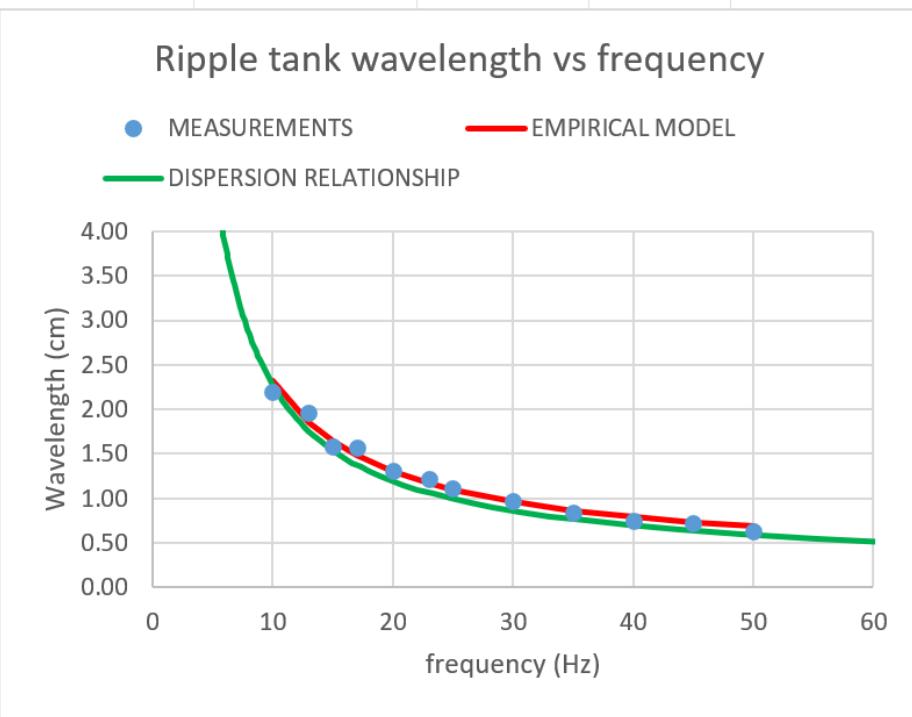
DISPERSION RELATIONSHIP

water density (kg/m³) 1000

lambda (cm)	k (m ⁻¹)	w ²	f (Hz)	1/f	phase velocity w/k (cm/s)
0.2	3.142E+03	2.296E+06	241.17	0.004	48.23
0.25	2.513E+03	1.184E+06	173.21	0.006	43.30
0.3	2.094E+03	6.917E+05	132.37	0.008	39.71
0.35	1.795E+03	4.403E+05	105.60	0.009	36.96
0.4	1.571E+03	2.986E+05	86.96	0.011	34.79
0.45	1.396E+03	2.126E+05	73.38	0.014	33.02
0.5	1.257E+03	1.573E+05	63.12	0.016	31.56
0.55	1.142E+03	1.201E+05	55.16	0.018	29.34

Dispersion relationship

$$\omega^2 = \left(\frac{\sigma k^3}{\rho_1} + gk \right) \tanh(kD)$$



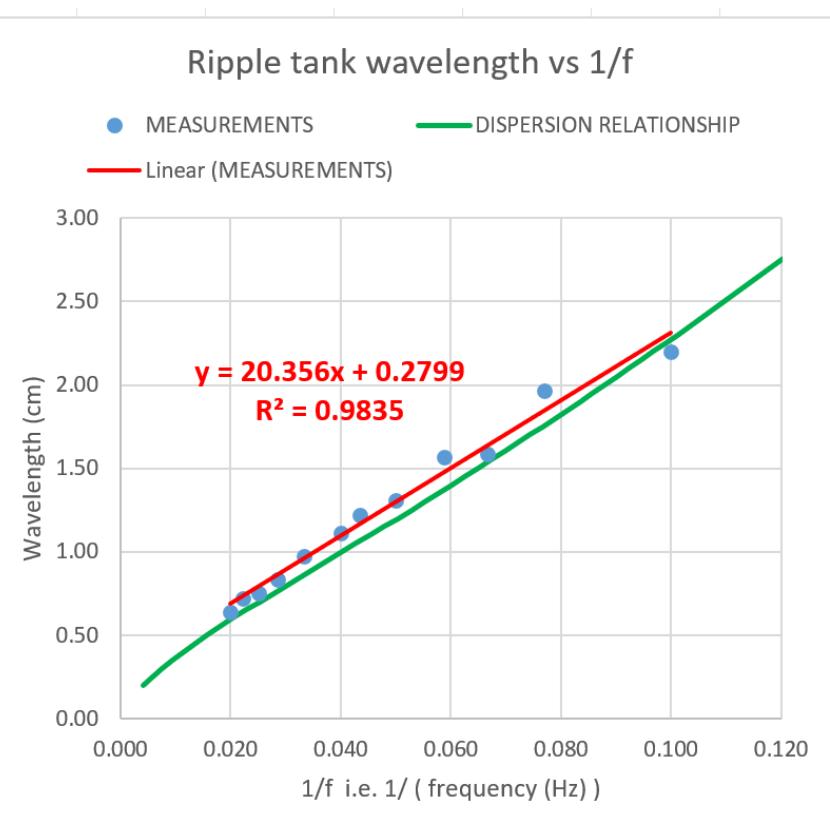
The dispersion relationship (green) is plotted against the measurements and an empirical model fit.

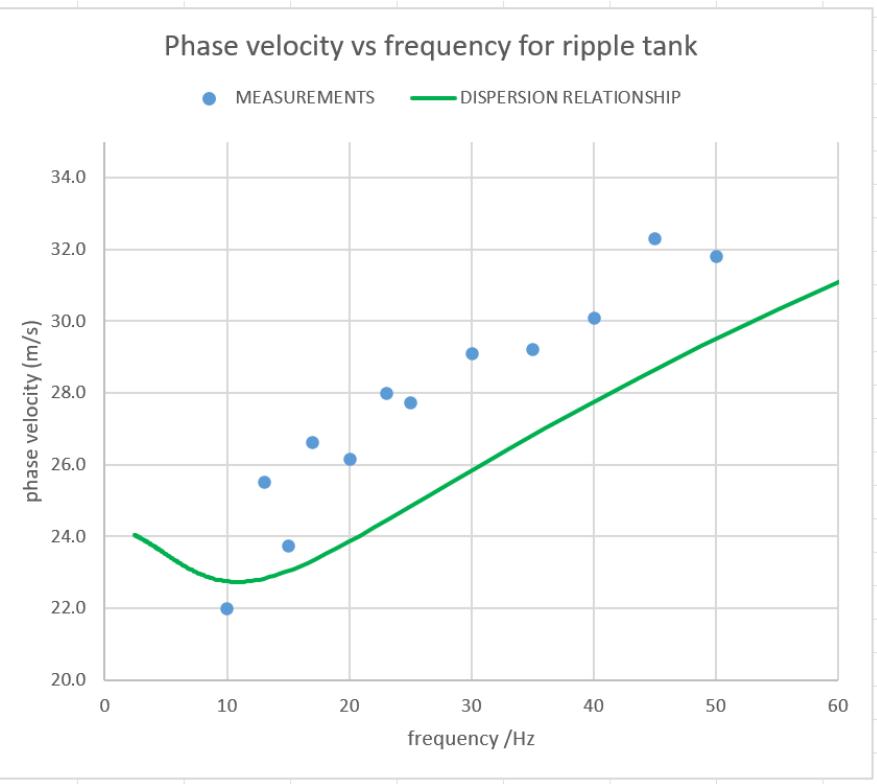
There is qualitative agreement (note error bars have been omitted in this initial analysis).

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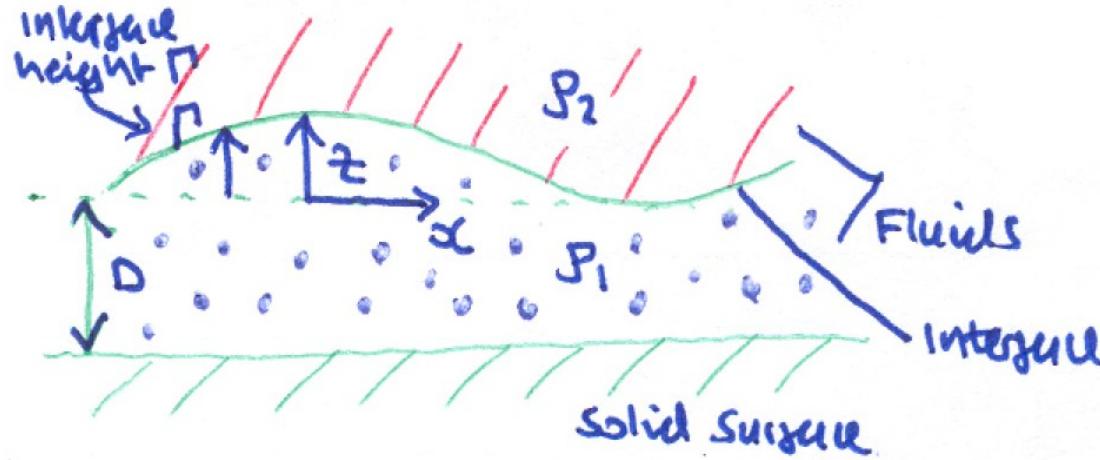
$$c_p = \frac{\omega}{k} = f\lambda, \quad c_g = \frac{d\omega}{dk}$$

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Phase velocity (measured) vs predictions from dispersion relationship is less well correlated, although perhaps follows the general upwards trend with the right gradient.
 Note there is a *minimum phase velocity for ripples*.

The measured phase velocity appears to be an overestimate, which is perhaps consistent with distortions caused by taking measurements from photographs.

It would be interesting to see a repeat of the experiment using only direct ruler-based measurements off the screen.



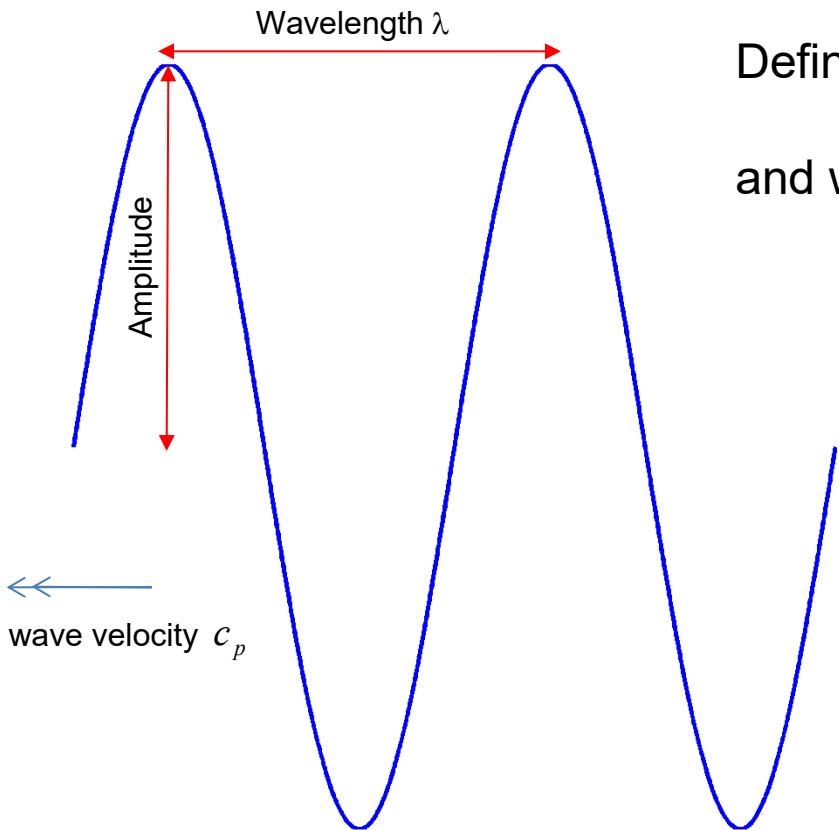
Mathematical theory of surface waves

Key topics:

Frequency
Wavenumber
Dispersion relationship

Phase and group velocity
Shallow and deep water waves
Minimum velocity of deep water ripples

A means of characterizing a wave-like disturbance is via a *dispersion relation*. This is an equation which relates the frequency of the wave to its wavelength, plus other parameters such as surface tension, fluid density, depth etc.



Define angular frequency

$$\omega = 2\pi f$$

and wavenumber

$$k = \frac{2\pi}{\lambda}$$

For deep water gravity waves, the dispersion relationship is

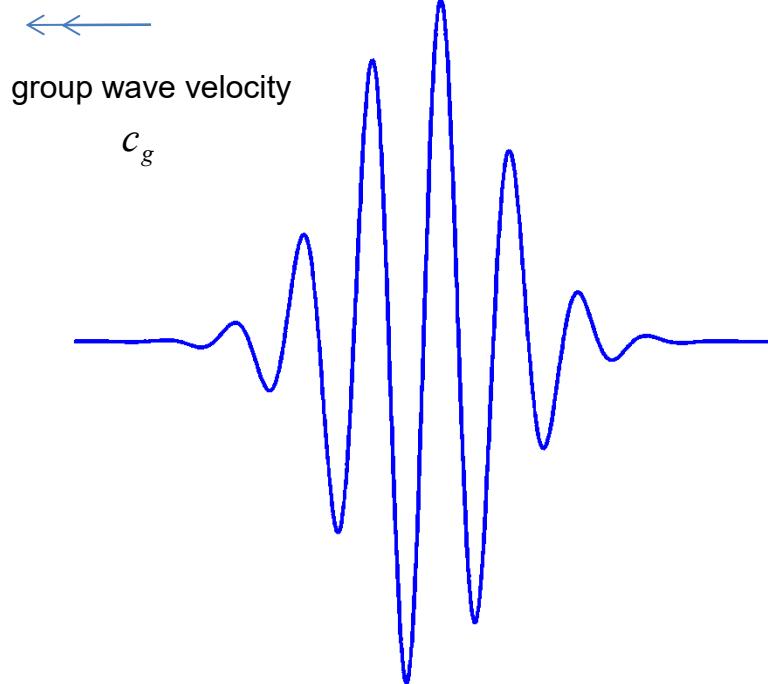
$$\omega^2 = gk$$

i.e. we ignore effects of the density of air, surface tension and assume the depth $D \gg \lambda$ i.e. $kD \gg 1$

The *phase velocity* (of individual wave crests) can be found from the dispersion relationship

$$c_p = f\lambda = \frac{\omega}{k}$$

The dispersion relationship allows to compute how fast groups of disturbances will travel. This *group velocity* can be *different* from the phase velocity. Relative motion of the wave crests to the overall envelope causes dispersion.



$$c_g = \frac{d\omega}{dk}$$

For deep water gravity driven waves

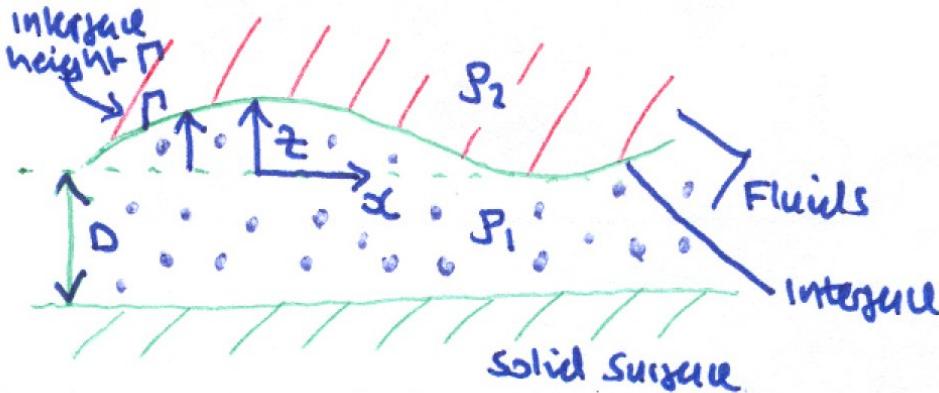
$$\omega = \sqrt{gk^{\frac{1}{2}}}$$

$$c_p = \frac{\omega}{k} = \sqrt{gk^{-\frac{1}{2}}}$$

$$c_g = \frac{d\omega}{dk} = \frac{1}{2} \sqrt{gk^{-\frac{1}{2}}}$$

$$\therefore c_g = \frac{1}{2} c_p$$

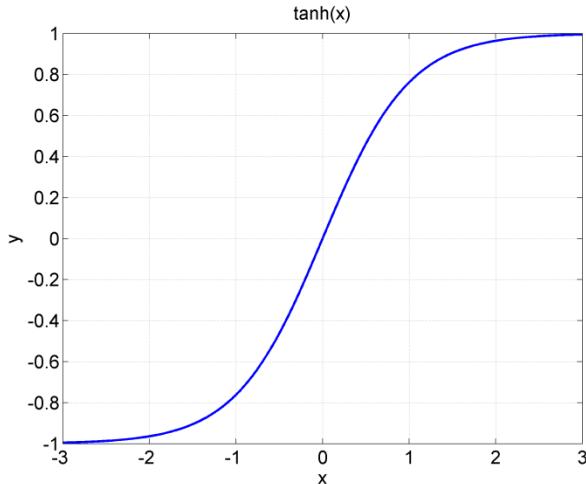
In general, for *vorticity* free waves on the interface of two *incompressible, Newtonian* fluids, the dispersion relationship can be shown to be, for waves with amplitude $\Gamma \ll D$



$$\omega^2 = \frac{\sigma k^3 + g(\rho_1 - \rho_2)k}{\rho_2 + \rho_1 \operatorname{cotanh}(kD)}$$

For waves on a water, air interface $\rho_1 \approx 1000\rho_2$, so ignore ρ_2

$$\omega^2 = \left(\frac{\sigma k^3}{\rho_1} + gk \right) \tanh(kD)$$



For shallow water waves $\tanh(kD) \approx kD$

$$\omega^2 = \frac{\sigma k^4 D}{\rho_1} + gk^2 D$$

For deep water waves $\tanh(kD) \approx 1$

$$\omega^2 = \frac{\sigma k^3}{\rho_1} + gk$$

Since $k = \frac{2\pi}{\lambda}$ the higher powers of k will contribute less for longer wavelengths

Hence for shallow water gravity waves $\omega^2 = gk^2 D$

e.g. waves coming ashore just before they break, or waves in a shallow river or canal

Similarly for **deep water waves** $\omega^2 = gk$

Note for deep water *ripples* (or 'capillary waves') this approximation is invalid.

We must use the full dispersion relation

$$\omega = \sqrt{\frac{\sigma k^3}{\rho_1} + gk}$$

Ripple phase velocity is:

$$c_p = \frac{\omega}{k} = \sqrt{\frac{\sigma k}{\rho_1} + \frac{g}{k}} = \sqrt{\frac{2\pi\sigma}{\lambda\rho_1} + \frac{g\lambda}{2\pi}}$$

Which has a *minima* at

$$c_p = \sqrt[4]{\frac{4g\sigma}{\rho_1}}$$

